

Curvature Tensors on Complex Lagrange Spaces

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Abstract. In this paper, our aim is to give an expression on $T_C(T'M)$ of the Levi-Civita connection for the Hermitian metric in a complex Lagrange space and to investigate its curvature tensor.

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1 Introduction

The study of the complex Lagrange geometry was recently developed by Gh. Munteanu ([6, 7, 8]). The holomorphic sectional curvature in a complex Finsler space was actually studied only on the horizontal part of Riemannian curvature ([1, 2]). In this paper, we begin a study of curvature tensors using the Sasaki lift on $T_C(T'M)$ of a general Lagrange metric tensor.

Let M be a complex manifold, $\dim_C M = n$, and let (U, z^k) be the complex coordinates in a local chart. The complexification $T_C M$ of the tangent bundle TM is decomposed in each $z \in M$ as $T_C M = T'M \oplus T''M$, where $T'M$ is the holomorphic bundle in which, as a complex manifold, a point is $u = (z^k, \eta^k)$ in a local chart. Let $\pi : T'M \rightarrow M$ be the canonical projection, $V(T'M) = \{\xi \in T(T'M) / \pi_*(\xi) = 0\}$ be the vertical subbundle, and $\mathcal{V}(T'M)$ the module of its vertical sections ([1, 2, 6]), spanned by $\{\dot{\partial} = \frac{\partial}{\partial \eta^k}\}$.

A complex nonlinear connection, (*c.n.c.*), is a supplementary subbundle to $V(T'M)$ in $T'(T'M)$, i.e.,

$$T'(T'M) = H(T'M) \oplus V(T'M) \tag{1.1}$$

and by $\mathcal{H}(T'M)$ we denote the horizontal distribution, in which an adapted base of the (*c.n.c.*) is $\{\delta_k = \frac{\delta}{\delta z^k} = \frac{\partial}{\partial z^k} - N_k^j \frac{\partial}{\partial \eta^j}\}$.

By conjugation we obtain a decomposition of whole $T_C(T'M)$:

$$T_C(T'M) = H(T'M) \oplus V(T'M) \oplus \overline{H(T'M)} \oplus \overline{V(T'M)} \tag{1.2}$$

and the corresponding conjugate bases will be denoted by $\{\delta_{\bar{k}} = \frac{\delta}{\delta \bar{z}^k}\}, \{\dot{\partial}_{\bar{k}} = \frac{\partial}{\partial \bar{\eta}^k}\}$.

2 The Levi-Civita connection on $T'M$

Let (M, L) be a complex Lagrange space, i.e., $L : T'M \rightarrow R$ is a smooth function for which the metric tensor $g_{i\bar{j}} = \partial^2 L / \partial \eta^i \partial \bar{\eta}^j$ is nondegenerate. We shall use through the Chern-Lagrange *c.n.c.*, introduced and studied by Gh. Munteanu ([8]):

$$N_j^i = g^{mi} \frac{\partial^2 L}{\partial z^j \partial \bar{\eta}^m} \quad (2.1)$$

in respect to which the corresponding adapted frames satisfy the following interesting relations

$$\begin{aligned} [\delta_j, \delta_k] = 0, \quad \delta_j(N_k^i) &= \delta_k(N_j^i), \quad (\partial_{\bar{i}} N_j^i) g_{i\bar{k}} - (\partial_{\bar{k}} N_j^i) g_{i\bar{l}} = 0 \\ (\partial_{\bar{i}} N_j^i) g_{i\bar{k}} + (\partial_{\bar{k}} N_j^i) g_{i\bar{l}} &= 2(\partial_{\bar{k}} g_{i\bar{l}}) N_j^i, \quad (\partial_{\bar{i}} N_j^i) g_{i\bar{k}} = \delta_j g_{i\bar{k}} \end{aligned} \quad (2.2)$$

In particular, when L is positively and absolutely homogeneous of degree two, $L(z, \lambda \eta) = |\lambda|^2 L(z, \eta)$, $\lambda \in \mathbb{C}$, then (M, L) is a complex Finsler space and $N_j^i = g^{mi} \frac{\partial g_{km}}{\partial z^j} \eta^k$ is just the Chern-Finsler connection ([1, 2]).

We consider now the Hermitian metrical structure G on $T_C(T'M)$, ([6]):

$$G = g_{i\bar{j}} dz^i \otimes d\bar{z}^j + g_{i\bar{j}} \delta \eta^i \otimes \delta \bar{\eta}^j \quad (2.3)$$

where $\{dz^k, \delta \eta^k\}$ is the adapted cobase of *c.n.c.*

First of all, we compute the coefficients of the Levi-Civita connection ∇ of in adapted frame $\{\delta_k, \partial_k, \delta_{\bar{k}}, \partial_{\bar{k}}\}$, following similar path as Anastasiei-Shimada did in the real case ([3]).

$$\begin{aligned} \nabla_{\delta_k} \delta_j &= L_{jk}^i \delta_i + A_{jk}^i \partial_i + A_{jk}^{\bar{i}} \delta_{\bar{i}} + A_{jk}^{\dot{i}} \dot{\partial}_i \\ \nabla_{\delta_k} \dot{\partial}_j &= B_{jk}^i \delta_i + L_{jk}^i \partial_i + B_{jk}^{\bar{i}} \delta_{\bar{i}} + B_{jk}^{\dot{i}} \dot{\partial}_i \\ \nabla_{\delta_k} \delta_{\bar{j}} &= D_{jk}^i \delta_i + D_{jk}^{\dot{i}} \dot{\partial}_i + L_{jk}^{\bar{i}} \delta_{\bar{i}} + D_{jk}^{\ddot{i}} \ddot{\partial}_i \\ \nabla_{\delta_k} \dot{\partial}_{\bar{j}} &= E_{jk}^i \delta_i + E_{jk}^{\dot{i}} \dot{\partial}_i + E_{jk}^{\bar{i}} \delta_{\bar{i}} + L_{jk}^{\dot{i}} \dot{\partial}_i \\ \nabla_{\partial_k} \delta_j &= C_{jk}^i \delta_i + F_{jk}^i \partial_i + F_{jk}^{\bar{i}} \delta_{\bar{i}} + F_{jk}^{\dot{i}} \dot{\partial}_i \\ \nabla_{\partial_k} \dot{\partial}_j &= G_{jk}^i \delta_i + C_{jk}^i \partial_i + G_{jk}^{\bar{i}} \delta_{\bar{i}} + C_{jk}^{\dot{i}} \dot{\partial}_i \\ \nabla_{\partial_k} \delta_{\bar{j}} &= H_{jk}^i \delta_i + H_{jk}^{\dot{i}} \dot{\partial}_i + C_{jk}^{\bar{i}} \delta_{\bar{i}} + H_{jk}^{\ddot{i}} \ddot{\partial}_i \\ \nabla_{\partial_k} \dot{\partial}_{\bar{j}} &= M_{jk}^i \delta_i + M_{jk}^{\dot{i}} \dot{\partial}_i + M_{jk}^{\bar{i}} \delta_{\bar{i}} + C_{jk}^{\dot{i}} \dot{\partial}_i \end{aligned} \quad (2.4)$$

and their conjugates, since $\nabla_{\bar{X}} \bar{Y} = \overline{\nabla_X Y}$.

Let T be the torsion of ∇ , i.e. $T(X, Y) = \nabla_X Y - \nabla_Y X - [X, Y]$, where X, Y are vector fields on $T_C(T'M)$. The condition ∇ is torsion-free is equivalent to

$$T(\delta_j, \delta_k) = T(\delta_j, \delta_{\bar{k}}) = T(\delta_j, \dot{\partial}_k) = T(\delta_{\bar{j}}, \dot{\partial}_k) = T(\dot{\partial}_j, \dot{\partial}_k) = T(\dot{\partial}_j, \dot{\partial}_{\bar{k}}) = 0 \quad (2.5)$$

and their conjugates. Using the following formulas

$$\begin{aligned} [\delta_j, \delta_k] &= R_{jk}^i \dot{\partial}_i, [\delta_j, \delta_{\bar{k}}] = (\delta_{\bar{k}} N_j^i) \dot{\partial}_i - (\delta_j \bar{N}_k^i) \dot{\partial}_{\bar{i}}, [\delta_j, \dot{\partial}_k] = (\dot{\partial}_k N_j^i) \dot{\partial}_i \\ [\delta_{\bar{j}}, \dot{\partial}_k] &= (\dot{\partial}_k \bar{N}_{\bar{j}}^i) \dot{\partial}_{\bar{i}}, [\dot{\partial}_j, \dot{\partial}_k] = 0, [\dot{\partial}_j, \dot{\partial}_{\bar{k}}] = 0, \end{aligned} \quad (2.6)$$

where $R_{jk}^i = \delta_k N_j^i - \delta_j N_k^i = 0$ for the Chern-Lagrange c.n.c., and their conjugates, by direct computation, requiring (2.5) and $\nabla G = 0$, we obtain

Theorem 2.1. *The local coefficients of the Levi-Civita connection of the Hermitian metrical structure G in the frame $\{\delta_i, \dot{\partial}_i, \delta_{\bar{i}}, \dot{\partial}_{\bar{i}}\}$ are as follows:*

$$\begin{aligned} L_{jk}^i &= \frac{1}{2} g^{i\bar{l}} (\delta_k g_{j\bar{l}} + \delta_j g_{k\bar{l}}), A_{jk}^i = A_{jk}^{\bar{i}} = A_{jk}^{\bar{i}} = 0 \\ B_{jk}^i &= \frac{1}{2} g^{i\bar{l}} (g_{j\bar{l}} \delta_k \bar{N}_l^h + \dot{\partial}_j g_{k\bar{l}}) = C_{kj}^i, L_{jk}^{\bar{i}} = g^{i\bar{l}} \delta_k g_{j\bar{l}}, B_{jk}^{\bar{i}} = B_{jk}^i = 0, \\ D_{jk}^i &= \frac{1}{2} g^{i\bar{l}} (\delta_j g_{k\bar{l}} + \delta_{\bar{l}} g_{k\bar{j}}) = L_{kj}^i, D_{jk}^{\bar{i}} = \frac{1}{2} g^{i\bar{l}} (g_{k\bar{l}} \delta_j \bar{N}_l^h - \dot{\partial}_{\bar{l}} g_{k\bar{j}}) \\ D_{jk}^{\bar{i}} &= \frac{1}{2} g^{i\bar{l}} (g_{l\bar{j}} \delta_k \bar{N}_j^h + \dot{\partial}_i g_{k\bar{j}}), F_{jk}^i = F_{jk}^{\bar{i}} = F_{jk}^{\bar{i}} = 0 \\ E_{jk}^i &= \frac{-1}{2} g^{i\bar{l}} (g_{k\bar{j}} \delta_{\bar{l}} \bar{N}_k^h - \dot{\partial}_j g_{k\bar{l}}) = C_{kj}^i, E_{jk}^{\bar{i}} = E_{jk}^i = L_{jk}^{\bar{i}} = 0 \\ G_{jk}^i &= g^{i\bar{l}} \dot{\partial}_k g_{j\bar{l}} \bar{N}_l^h, C_{jk}^i = g^{i\bar{l}} \dot{\partial}_k g_{j\bar{l}}, G_{jk}^{\bar{i}} = G_{jk}^i = 0 \\ H_{jk}^i &= -g^{i\bar{l}} \dot{\partial}_k g_{l\bar{j}} \bar{N}_j^h, H_{jk}^{\bar{i}} = H_{jk}^i = M_{jk}^i = M_{jk}^{\bar{i}} = M_{jk}^{\bar{i}} = C_{jk}^i = 0 \end{aligned} \quad (2.7)$$

and the conjugates.

Remark 2.1. In particular, if $\delta_k g_{j\bar{l}} = \delta_j g_{k\bar{l}}$, i.e. the metric tensor is strongly Kähler, ([1]), then $L_{jk}^i = L_{jk}^{\bar{i}}$, and $D_{jk}^i = D_{jk}^{\bar{i}} = 0$.

3 The curvature tensor of the Levi-Civita connection on $T'M$

As in real case, ([3]), we shall compute the components of the curvatures of ∇ using an intermediate arbitrary d -connection, D on $T_C(T'M)$ ([6]):

$$\begin{aligned} D_{\delta_k} \delta_j &= L_{jk}^i \delta_i, D_{\delta_k} \dot{\partial}_j = L_{jk}^i \dot{\partial}_i, D_{\delta_k} \delta_{\bar{j}} = L_{jk}^{\bar{i}} \delta_{\bar{i}}, D_{\delta_k} \dot{\partial}_{\bar{j}} = L_{jk}^{\bar{i}} \dot{\partial}_{\bar{i}} \\ D_{\dot{\partial}_k} \delta_j &= C_{jk}^i \delta_i, D_{\dot{\partial}_k} \dot{\partial}_j = C_{jk}^i \dot{\partial}_i, D_{\dot{\partial}_k} \delta_{\bar{j}} = C_{jk}^{\bar{i}} \delta_{\bar{i}}, D_{\dot{\partial}_k} \dot{\partial}_{\bar{j}} = C_{jk}^{\bar{i}} \dot{\partial}_{\bar{i}} \end{aligned} \quad (3.1)$$

and their conjugates.

Let us express the curvature components of D , which is not free torsion, in adapted base.

$$\begin{aligned}
R(\delta_h, \delta_k)\delta_j &= \tilde{R}_{jkh}^i \delta_i, & R(\delta_h, \delta_k)\delta_{\bar{j}} &= \tilde{R}_{\bar{j}kh}^{\bar{i}} \delta_{\bar{i}}, & R(\delta_h, \delta_k)\delta_{\bar{j}} &= \tilde{R}_{\bar{j}kh}^{\bar{i}} \delta_{\bar{i}} \\
R(\delta_h, \delta_k)\dot{\delta}_j &= \tilde{\Omega}_{jkh}^i \dot{\delta}_i, & R(\delta_h, \delta_k)\dot{\delta}_{\bar{j}} &= \tilde{\Omega}_{\bar{j}kh}^{\bar{i}} \dot{\delta}_{\bar{i}}, & R(\delta_h, \delta_k)\dot{\delta}_{\bar{j}} &= \tilde{\Omega}_{\bar{j}kh}^{\bar{i}} \dot{\delta}_{\bar{i}} \\
R(\delta_h, \dot{\delta}_k)\delta_j &= \tilde{\Pi}_{jkh}^i \delta_i, & R(\delta_h, \dot{\delta}_k)\delta_{\bar{j}} &= \tilde{\Pi}_{\bar{j}kh}^{\bar{i}} \delta_{\bar{i}}, & R(\delta_h, \dot{\delta}_k)\delta_{\bar{j}} &= \tilde{\Pi}_{\bar{j}kh}^{\bar{i}} \delta_{\bar{i}} \\
R(\delta_h, \dot{\delta}_k)\dot{\delta}_j &= \tilde{P}_{jkh}^i \dot{\delta}_i, & R(\delta_h, \dot{\delta}_k)\dot{\delta}_{\bar{j}} &= \tilde{P}_{\bar{j}kh}^{\bar{i}} \dot{\delta}_{\bar{i}}, & R(\delta_h, \dot{\delta}_k)\dot{\delta}_{\bar{j}} &= \tilde{P}_{\bar{j}kh}^{\bar{i}} \dot{\delta}_{\bar{i}} \\
R(\delta_h, \dot{\delta}_{\bar{k}})\delta_j &= \tilde{\Theta}_{jkh}^i \delta_i, & R(\delta_h, \dot{\delta}_{\bar{k}})\delta_{\bar{j}} &= \tilde{Q}_{\bar{j}kh}^{\bar{i}} \delta_{\bar{i}} \\
R(\dot{\delta}_h, \dot{\delta}_k)\delta_j &= \tilde{\Xi}_{jkh}^i \delta_i, & R(\dot{\delta}_h, \dot{\delta}_k)\delta_{\bar{j}} &= \tilde{\Xi}_{\bar{j}kh}^{\bar{i}} \delta_{\bar{i}}, & R(\dot{\delta}_h, \dot{\delta}_k)\delta_{\bar{j}} &= \tilde{\Xi}_{\bar{j}kh}^{\bar{i}} \delta_{\bar{i}} \\
R(\dot{\delta}_h, \dot{\delta}_k)\dot{\delta}_j &= \tilde{S}_{jkh}^i \dot{\delta}_i, & R(\dot{\delta}_h, \dot{\delta}_k)\dot{\delta}_{\bar{j}} &= \tilde{S}_{\bar{j}kh}^{\bar{i}} \dot{\delta}_{\bar{i}}, & R(\dot{\delta}_h, \dot{\delta}_k)\dot{\delta}_{\bar{j}} &= \tilde{S}_{\bar{j}kh}^{\bar{i}} \dot{\delta}_{\bar{i}}
\end{aligned} \tag{3.2}$$

and their conjugates, because $R(X, Y) = -R(Y, X)$.

The local components of curvature tensor R can be obtained easily from (3.1) and (3.2).

Our interest is to compute the curvature tensor K of ∇ by means of R . First, let us decompose the Riemannian curvature K in respect to the adapted base of Chern-Lagrange *c.n.c.*

$$K(\delta_h, \delta_k)\delta_j = R_{jkh}^h \delta_i + R_{jkh}^v \dot{\delta}_i + R_{jkh}^{\bar{h}} \delta_{\bar{i}} + R_{jkh}^v \dot{\delta}_{\bar{i}} \tag{3.3}$$

and similar for $K(\delta_h, \delta_k)\delta_{\bar{j}}$, $K(\delta_h, \delta_k)\delta_{\bar{j}}$, and then for $K(\delta_h, \delta_k)\dot{\delta}_j$ with coefficients Ω_{jk}^h , etc., at last being $K(\dot{\delta}_h, \dot{\delta}_k)\dot{\delta}_j$ with coefficients S_{jk}^h , and their conjugates.

Now, we can give directly an expression of Riemannian curvature K as a function of curvature R .

Theorem 3.1. *The local components of curvature of Levi-Civita connection of G metric on $T_C(T^*M)$ in respect to the adapted base of Chern-Lagrange *c.n.c.* are:*

$$\begin{aligned}
R_{jkh}^h &= \tilde{R}_{jkh}^i, & R_{jkh}^v &= R_{jkh}^{\bar{h}} = R_{jkh}^{\bar{v}} = 0 \\
R_{jkh}^i &= A_{hk} \{ L_{h\bar{j}|k}^3 + D_{\bar{j}k}^2 C_{hl}^1 + D_{\bar{j}k}^4 C_{h\bar{l}}^3 \}, & R_{jkh}^v &= A_{hk} \{ D_{\bar{j}h|k}^2 \} \\
R_{jkh}^{\bar{h}} &= \tilde{R}_{jkh}^{\bar{i}}, & R_{jkh}^{\bar{v}} &= A_{hk} \{ D_{\bar{j}h|k}^4 \}, & R_{jkh}^h &= -\{ L_{h\bar{j}|\bar{k}}^3 + L_{h\bar{k}|\bar{j}}^3 \} \\
R_{jkh}^{\bar{i}} &= \tilde{R}_{\bar{j}kh}^{\bar{i}} - L_{h\bar{j}}^3 L_{\bar{k}l}^3 - C_{\bar{k}l}^3 D_{\bar{j}h}^2 - C_{\bar{k}l}^1 D_{\bar{j}h}^4 \\
R_{jkh}^{\bar{v}} &= -D_{\bar{j}h|k}^4 - L_{h\bar{k}}^3 D_{\bar{j}l}^4 - L_{h\bar{j}}^3 D_{\bar{k}l}^2 - (\delta_{\bar{k}}^l N_{\bar{h}}^l) H_{\bar{j}l}^{CL} \\
\Omega_{jkh}^h &= A_{hk} \{ C_{k\bar{j}|h}^1 \}, & \Omega_{jkh}^v &= \Omega_{jkh}^{\bar{h}} = \Omega_{jkh}^{\bar{v}} = 0 \\
\Omega_{jkh}^i &= A_{hk} \{ C_{h\bar{j}|k}^3 + L_{\bar{j}k}^3 C_{h\bar{l}}^3 \}, & \Omega_{jkh}^v &= \Omega_{jkh}^{\bar{h}} = \Omega_{jkh}^{\bar{v}} = 0
\end{aligned}$$

$$\begin{aligned}
 \Omega_{j\bar{k}h}^h &= -C_{h\bar{j}}^3 - L_{j\bar{k}}^3 C_{l\bar{j}}^3 + C_{k\bar{j}}^1 L_{h\bar{l}}^3, \quad \Omega_{j\bar{k}h}^v = C_{h\bar{j}}^1 D_{l\bar{h}}^2 - C_{h\bar{j}}^3 D_{l\bar{k}}^4 \\
 \Omega_{j\bar{k}h}^{\bar{h}} &= C_{k\bar{j}}^1 \bar{L}_{h\bar{l}}^3 + L_{k\bar{h}}^3 C_{l\bar{j}}^1 - C_{h\bar{j}}^3 L_{h\bar{l}}^3 + (\delta_h N_k^l) C_{j\bar{l}}^1 \\
 \Omega_{j\bar{k}h}^{\bar{v}} &= \tilde{\Omega}_{j\bar{k}h}^{\bar{v}} + C_{k\bar{j}}^1 D_{l\bar{h}}^4 - C_{h\bar{j}}^3 D_{l\bar{k}}^2 \\
 \Pi_{j\bar{k}h}^h &= \tilde{\Pi}_{j\bar{k}h}^h, \quad \Pi_{j\bar{k}h}^v = \Pi_{j\bar{k}h}^{\bar{h}} = \Pi_{j\bar{k}h}^{\bar{v}} = 0, \quad \Pi_{j\bar{k}h}^{\bar{v}} = -D_{j\bar{h}|k}^2 - C_{h\bar{k}}^1 D_{j\bar{l}}^2 \\
 \Pi_{j\bar{k}h}^{\bar{h}} &= -L_{h\bar{j}}^3 - C_{h\bar{k}}^1 L_{l\bar{j}}^3 + H_{j\bar{k}}^4 C_{h\bar{l}}^1 - D_{j\bar{h}}^2 C_{l\bar{h}}^1, \\
 \Pi_{j\bar{k}h}^{\bar{v}} &= \tilde{\Pi}_{j\bar{k}h}^{\bar{v}}, \quad \tilde{\Pi}_{j\bar{k}h}^{\bar{v}} = C_{j\bar{k}}^3 D_{l\bar{h}}^4 - L_{j\bar{h}}^3 H_{l\bar{k}}^4 - (\partial_k N_h^l) H_{j\bar{l}}^1, \\
 \Pi_{j\bar{k}h}^h &= -L_{h\bar{j}}^3 - C_{j\bar{k}}^1 L_{l\bar{j}}^3, \quad \Pi_{j\bar{k}h}^v = -D_{j\bar{h}|k}^2 - C_{h\bar{k}}^3 D_{j\bar{l}}^2 - L_{h\bar{j}}^3 \Pi_{l\bar{k}}^4, \\
 \Pi_{j\bar{k}h}^{\bar{h}} &= \tilde{\Pi}_{j\bar{k}h}^{\bar{h}} - D_{j\bar{h}}^4 C_{l\bar{k}h}^1, \quad \Pi_{j\bar{k}h}^{\bar{v}} = -D_{j\bar{h}|k}^4 - C_{h\bar{k}}^3 D_{j\bar{l}}^4 - (\partial_k N_h^l) H_{j\bar{l}}^4, \\
 P_{j\bar{k}h}^h &= G_{j\bar{k}|h}^1 - C_{h\bar{j}}^1 C_{l\bar{k}}^1 - C_{h\bar{k}}^1 B_{j\bar{l}}^1, \quad P_{j\bar{k}h}^v = \tilde{P}_{j\bar{k}h}^v, \quad P_{j\bar{k}h}^{\bar{h}} = P_{j\bar{k}h}^{\bar{v}} = 0, \\
 P_{j\bar{k}h}^{\bar{h}} &= -C_{h\bar{j}}^3 - C_{k\bar{h}}^1 C_{l\bar{j}}^3, \quad P_{j\bar{k}h}^v = P_{j\bar{k}h}^{\bar{h}} = P_{j\bar{k}h}^{\bar{v}} = 0, \\
 P_{j\bar{k}h}^{\bar{v}} &= -C_{h\bar{j}}^3 - C_{h\bar{k}}^3 C_{l\bar{j}}^3 + G_{j\bar{k}}^1 L_{h\bar{l}}^3, \quad P_{j\bar{k}h}^v = G_{j\bar{k}}^1 D_{l\bar{h}}^2 - C_{h\bar{j}}^3 H_{l\bar{k}}^4, \\
 P_{j\bar{k}h}^h &= G_{j\bar{k}|h}^1, \quad P_{j\bar{k}h}^{\bar{v}} = \tilde{P}_{j\bar{k}h}^{\bar{v}} + C_{j\bar{k}}^1 D_{l\bar{h}}^4, \quad \Theta_{j\bar{k}h}^h = \tilde{\Theta}_{j\bar{k}h}^h + H_{j\bar{k}}^4 C_{h\bar{l}}^1, \\
 \Theta_{j\bar{k}h}^v &= H_{j\bar{k}|h}^4, \quad \Theta_{j\bar{k}h}^{\bar{h}} = \Theta_{j\bar{k}h}^{\bar{v}} = 0, \\
 Q_{j\bar{k}h}^h &= -C_{h\bar{j}}^1 - C_{h\bar{k}}^3 C_{l\bar{j}}^1 - (\partial_k N_h^l) C_{j\bar{l}}^1, \quad Q_{j\bar{k}h}^v = -\tilde{Q}_{j\bar{k}h}^v - C_{h\bar{j}}^1 H_{l\bar{k}}^4, \\
 Q_{j\bar{k}h}^{\bar{h}} &= Q_{j\bar{k}h}^{\bar{v}} = 0, \quad \Xi_{j\bar{k}h}^h = \tilde{\Xi}_{j\bar{k}h}^h, \quad \Xi_{j\bar{k}h}^v = \Xi_{j\bar{k}h}^{\bar{h}} = \Xi_{j\bar{k}h}^{\bar{v}} = 0, \\
 \Xi_{j\bar{k}h}^{\bar{h}} &= \Xi_{j\bar{k}h}^{\bar{v}} = 0, \quad \tilde{\Xi}_{j\bar{k}h}^{\bar{h}} = \tilde{\Xi}_{j\bar{k}h}^{\bar{h}}, \quad \tilde{\Xi}_{j\bar{k}h}^{\bar{v}} = A_{kh} \{H_{j\bar{k}|h}^4\}, \quad \Xi_{j\bar{k}h}^h = \Xi_{j\bar{k}h}^v = 0, \\
 H_{j\bar{k}h}^{\bar{h}} &= \tilde{H}_{j\bar{k}h}^{\bar{h}} - H_{j\bar{h}}^4 C_{l\bar{k}}^1, \quad H_{j\bar{k}h}^{\bar{v}} = -H_{j\bar{h}|k}^4, \quad S_{j\bar{k}h}^h = A_{kh} \{G_{j\bar{k}|h}^1\}, \\
 S_{j\bar{k}h}^v &= S_{j\bar{k}h}^{\bar{h}} = S_{j\bar{k}h}^{\bar{v}} = 0, \quad S_{j\bar{k}h}^h = S_{j\bar{k}h}^v = S_{j\bar{k}h}^{\bar{h}} = S_{j\bar{k}h}^{\bar{v}} = 0, \quad S_{j\bar{k}h}^h = S_{j\bar{k}h}^v = 0, \\
 S_{j\bar{k}h}^{\bar{h}} &= G_{j\bar{k}|h}^1, \quad S_{j\bar{k}h}^{\bar{v}} = \tilde{S}_{j\bar{k}h}^{\bar{v}} + G_{j\bar{k}}^1 H_{l\bar{h}}^4,
 \end{aligned}$$

and their conjugates, where by $|k, ||k, \bar{|}k, \bar{||}k$ we denote $h-, v-, \bar{h}, \bar{v}-$ covariant derivatives with respect to D , and $A_{hk} \{T_{hk}^i\} = T_{hk}^i - T_{kh}^i$.

Remark 3.1. If $\delta_k g_{j\bar{i}} = \delta_j g_{k\bar{i}}$, then we shall obtain the particular form of Theorem 3.2.

Although the paper has some technical elements, this calculus is absolutely necessary for a further study of holomorphic sectional curvature.

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