

AN ASYMPTOTIC EXPANSION FOR SOME SECOND ORDER DIFFERENTIAL EQUATIONS IN HILBERT SPACES

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Abstract. An asymptotic expansion (series) for the solution u_ε of the boundary value problem (A_ε) below is constructed. The n -th partial sum of this series is an asymptotic approximation for u_ε in $[0, T]$ to within accuracy of the order $\varepsilon^{n+1/2}$. In particular, u_ε approaches the solution u of the first order problem (A_0) , in the entire interval $[0, T]$, if $b = u(T)$ and in $[0, \delta]$, for some positive δ , if $b \neq u(T)$.

1. INTRODUCTION

In this paper, we intend to compare the solution of the boundary value problem:

$$\begin{cases} \varepsilon u_\varepsilon'' - u_\varepsilon' = Au_\varepsilon + f & \text{a.e. on } [0, T] \\ u_\varepsilon(0) = a, \quad u_\varepsilon(T) = b, \end{cases} \quad (A_\varepsilon)$$

with the solution of the problem

$$\begin{cases} u' + Au = -f, & \text{a.e. on } [0, T] \\ u(0) = a. \end{cases} \quad (A_0)$$

Here ε is a small positive parameter, A is a linear operator on a Hilbert space H , f is a given function and $a, b \in D(A)$ are given. If in (A_ε) we take $\varepsilon = 0$, we obtain the problem (A_0) , so (A_ε) is called the perturbed problem and (A_0) is the unperturbed problem.

Our main goal is to show that the difference $u_\varepsilon - u$ (in a suitably chosen measure) approximates zero on a subset $[0, \delta]$, for some δ , with $0 < \delta < T$. Moreover, since $u(t)$ does not approximate $u_\varepsilon(t)$ for small ε in the so-called boundary layer $(\delta, T]$, we shall construct an asymptotic expansion (or asymptotic series) for $u_\varepsilon(t)$, that is valid in the entire interval $[0, T]$.

First, let us establish the hypotheses under which we work and recall some existence results concerning the solutions of (A_ε) and (A_0) .

We are given a real Hilbert space H (with the norm $\|\cdot\|$) and an operator $A: H \rightarrow H$ such that

(I1) A is a single-valued, linear, densely defined, symmetric and positive definite operator, i.e., there is a positive number ω , such that

$$(Ax, x) \geq \omega \|x\|^2, \quad (\forall)x \in H. \quad (1.1)$$

$$(12) \quad a, b \in D(A) \quad (1.2)$$

$$(13) \quad f \in W^{1,2}(0, T; H), \quad \text{i.e. } f, f' \in L^2(0, T; H). \quad (1.3)$$

We know that, if A is maximal monotone set in $H \times H$, $a \in D(A)$ and $f \in W^{1,1}(0, T; H)$, then (A_0) has a unique strong solution u , which belongs to $W^{1,\infty}(0, T; H)$ ([6], [8], [9], [10]), i.e. $u(t)$ is absolutely continuous on $[0, T]$ and $du/dt \in L^\infty(0, T; H)$.

The problem A_ε is a particular case of the more general boundary value problem

$$\begin{cases} pu'' + ru' \in Au + f, & \text{a.e. on } [0, T] \\ u'(0) \in \alpha(u(0) - a), & u'(T) \in -\beta(u(T) - b), \end{cases} \quad (1.4)$$

where A, α, β are maximal monotone operators in a real Hilbert space H , with $0 \in D(A) \cap D(\alpha) \cap D(\beta)$, $0 \in \alpha(0)$, $0 \in \beta(0)$, $a, b \in D(A)$, $f \in L^2(0, T; H)$ and $p, r : [0, T] \rightarrow \mathbb{R}$, $p, r \in W^{1,\infty}(0, T)$. One also assumes that, if A_λ is the Yosida approximation of A , i.e. $A_\lambda x = (1/\lambda)(x - (I + \lambda A)^{-1}x)$, then

$$\begin{aligned} (A_\lambda x - A_\lambda y, z) &\geq 0, \quad (\forall)z \in \alpha(x - y), \quad \text{where } x - y \in D(\alpha), \\ (A_\lambda x - A_\lambda y, w) &\leq 0, \quad (\forall)w \in -\beta(x - y), \quad \text{where } x - y \in D(\beta). \end{aligned}$$

The existence of the solutions of (1.4) was studied by L. Véron [12] and by R. Aftabizadeh and N.H. Pavel [2], [3]. One of these results ([2]) establishes that, if A is maximal strongly monotone operator in H , then (1.4) has a unique solution $u \in W^{2,2}(0, T; H)$.

Similar results for second order differential equations on half-axis, were obtained by N. Apreutesei [4], [5]. We recall one of these theorems, which we need later. Consider the problem

$$\begin{cases} pu'' + ru' \in Bu + g, & \text{a.e. on } [0, \infty) \\ u'(0) \in \alpha(u(0) - a), \end{cases} \quad (1.5)$$

where B and α are maximal monotone operators in H , with $0 \in D(B) \cap D(\alpha)$, $0 \in \alpha(0)$, $a \in D(B)$ and $p, r \in W^{1,\infty}(0, \infty)$, $p(t) \geq c > 0$, $r(t) \geq r_0 > 0$, $(\forall)t \in [0, \infty)$. Let \tilde{r} be the function given by

$$\tilde{r}(t) = \exp\left(\int_0^t (r(s)/p(s))ds\right) \quad (1.6)$$

and $L^2_{\tilde{r}}(0, \infty; H)$ be the space $L^2(0, \infty; H)$ with the weight function \tilde{r} , i.e. the inner product in $L^2_{\tilde{r}}(0, \infty; H)$ is

$$\langle\langle u, v \rangle\rangle = \int_0^\infty \tilde{r}(t)(u(t), v(t))dt, \quad (\forall)u, v \in L^2_{\tilde{r}}(0, \infty; H). \quad (1.7)$$

THEOREM 1.1. *If the above hypotheses hold, $D(\alpha)$ is bounded in H , $f \in L^2_{\tilde{r}}(0, \infty; H)$ and*

$$(A_\lambda x - A_\lambda y, z) \geq 0, \quad \forall z \in \alpha(x - y), \quad \text{where } x - y \in D(\alpha), \quad (1.8)$$

then (1.5) has at least one solution $u \in H^2_{\tilde{r}}(0, \infty; H) = \{u/u, u', u'' \in L^2_{\tilde{r}}(0, \infty; H)\}$.

We come back to (1.4) and put $\alpha = \beta = \partial j$, where $j : H \rightarrow (-\infty, +\infty]$,

$$j(x) = \begin{cases} 0, & x = 0 \\ +\infty, & \text{otherwise.} \end{cases} \quad (1.9)$$

Then $D(\partial j) = \{0\}$ and $\partial j(0) = H$, so the boundary conditions become $u(0) = a$, $u(T) = b$, as in (A_ε) . Taking $p(t) = \varepsilon$, $r(t) = -1$, $(\forall)t \in [0, T]$, we obtain the problem (A_ε) .

Therefore, under the hypotheses (I1)-(I3), the problem (A_ε) has a unique solution $u_\varepsilon \in W^{2,2}(0, T; H)$. Moreover, as we shall prove in section 2, if $f \in W^{1,2}(0, T; H)$, then $u_\varepsilon \in W^{2,\infty}(0, T; H)$ (see Theorem 2.1).

The motivation of this paper is:

1. First of all, the elliptic equation from (1.4) is too abstract. If we particularise $p \equiv \varepsilon$, $r \equiv -1$ and A, α, β as above, we obtain the problem (A_ε) . It's solution, u_ε , converges to the solution u of (A_0) and the problem (A_0) is the model of many physical phenomenon. Therefore, (A_ε) can be interpreted as an approximation (in a certain sense) of the mathematical models of these processes or phenomenon.
2. On the other hand, we shall prove that, under the assumptions (II)-(I3), u_ε is in $W^{2,\infty}(0, T; H)$, so the solution $u \in W^{1,\infty}(0, T; H)$ of (A_0) is approximated by the smoother function u_ε . This justifies the necessity of the elliptic regularization.
3. We show that, if $b = u(T)$, then $u_\varepsilon - u \rightarrow 0$, as $\varepsilon \rightarrow 0$, in $C(0, T; H)$ and, moreover this difference is of order ε to a positive power; if $b \neq u(T)$, then u_ε approximates u just outside of a neighbourhood $(\delta, T]$ of T .
4. Moreover, we construct an approximation with higher accuracy and the complete asymptotic expansion for the solution of the problem (A_ε) , in the spirit of A. Vasilieva, V. Butuzov and L. Kalachev [11].

Other papers which deal with the approximation of the solutions of an unperturbed problem by the solutions of a perturbed problem are [1], [7].

The structure of this work is the following. In the second section, we give a regularization result for the problem (A_ε) . The next paragraph is concerned with the construction of the asymptotic expansion (series) of the solution u_ε of the perturbed problem (A_ε) . In section 4, we justify the asymptotic expansion and we estimate the remainder term. In the last paragraph, an application of this theory to a partial differential equation is given.

2. A REGULARITY RESULT

As we mention above, the solution u_ε of (A_ε) is in $W^{2,2}(0, T; H)$. We shall establish that, in fact, u_ε belongs to $W^{2,\infty}(0, T; H)$, which explain why we approximate the solution $u \in W^{1,\infty}(0, T; H)$ of (A_0) by the solution u_ε of (A_ε) . More exactly, we give

THEOREM 2.1. *If $f \in W^{1,2}(0, T; H)$ and A is a linear operator satisfying condition (II), the solution u_ε of (A_ε) is in $W^{2,\infty}(0, T; H)$ and, for every $\varepsilon > 0$, we have the estimate*

$$\begin{aligned} \frac{\varepsilon}{2} \|u_\varepsilon''(t)\|^2 &\leq \frac{1}{\varepsilon} \|u_\varepsilon'(0)\|^2 + \frac{2}{\varepsilon} \|Aa\|^2 + \frac{2}{\varepsilon} \|f(0)\| + 2 \int_0^t \|u_\varepsilon''(s)\|^2 ds + \\ &+ \int_0^t \|f'(s)\|^2 ds + \frac{1}{2} \sup_{t \in [0, T]} (Au_\varepsilon'(t), u_\varepsilon'(t)), \quad t \in [0, T]. \end{aligned} \tag{2.1}$$

Proof. We know that, for every fixed positive ε , $u_\varepsilon \in W^{2,2}(0, T; H)$, so $u_\varepsilon, u_\varepsilon' \in L^\infty(0, T; H)$. Without loss of generality, suppose that A is a bounded operator. If not, we replace A by A_λ , we denote by $u_{\varepsilon\lambda}$ the solution of the corresponding approximation problem and we take the limit for $\lambda \downarrow 0$, taking into account the fact that $u_{\varepsilon\lambda} \rightarrow u_\varepsilon$, $u'_{\varepsilon\lambda} \rightarrow u'_\varepsilon$, $u''_{\varepsilon\lambda} \rightarrow u''_\varepsilon$ and $A_\lambda u_{\varepsilon\lambda} \rightarrow Au_\varepsilon$; thus we obtain (2.1). Let us show now that $u''_\varepsilon \in L^\infty(0, T; H)$ and (2.1).

To do this, we write the equation in $t+h$ and in t , we subtract them and we multiply

this difference by $u'_\varepsilon(t+h) - u'_\varepsilon(t)$. One obtains

$$\begin{aligned} \varepsilon(u''_\varepsilon(t+h) - u''_\varepsilon(t), u'_\varepsilon(t+h) - u'_\varepsilon(t)) - \|u'_\varepsilon(t+h) - u'_\varepsilon(t)\|^2 = \\ (Au_\varepsilon(t+h) - Au_\varepsilon(t), u'_\varepsilon(t+h) - u'_\varepsilon(t)) + \\ + (f(t+h) - f(t), u'_\varepsilon(t+h) - u'_\varepsilon(t)), \text{ a.e. on } (0, T-h). \end{aligned} \quad (2.2)$$

One integrates (2.2) on $[0, t]$ (with $0 \leq t \leq T$), one uses the linearity of A and the fact that, for $v(t) = u_\varepsilon(t+h) - u_\varepsilon(t)$,

$$\int_0^t (Av, v') ds = \frac{1}{2}(Av(t), v(t)) - \frac{1}{2}(Av(0), v(0)). \quad (2.3)$$

We deduce that

$$\begin{aligned} \frac{\varepsilon}{2} \|u'_\varepsilon(t+h) - u'_\varepsilon(t)\|^2 - \frac{\varepsilon}{2} \|u'_\varepsilon(h) - u'_\varepsilon(0)\|^2 - \int_0^t \|u'_\varepsilon(s+h) - u'_\varepsilon(s)\|^2 ds \leq \\ \leq \frac{1}{2} (A(u_\varepsilon(t+h) - u_\varepsilon(t)), u_\varepsilon(t+h) - u_\varepsilon(t)) + \\ + \int_0^t (f(s+h) - f(s), u'_\varepsilon(s+h) - u'_\varepsilon(s)) ds, \text{ a.e. on } (0, T-h). \end{aligned} \quad (2.4)$$

Let us estimate now $\|u'_\varepsilon(h) - u'_\varepsilon(0)\|^2$. Integrating the given equation on $[0, h]$, we find

$$\varepsilon \|u'_\varepsilon(h) - u'_\varepsilon(0)\| \leq \|u_\varepsilon(h) - u_\varepsilon(0)\| + \int_0^h \|Au_\varepsilon(s)\| ds + \int_0^h \|f(s)\| ds,$$

so

$$\frac{\varepsilon}{2} \|u'_\varepsilon(h) - u'_\varepsilon(0)\|^2 \leq \frac{1}{\varepsilon} \|u_\varepsilon(h) - u_\varepsilon(0)\|^2 + \frac{2}{\varepsilon} \left(\int_0^h \|Au_\varepsilon(s)\| ds \right)^2 + \frac{2}{\varepsilon} \left(\int_0^h \|f(s)\| ds \right)^2. \quad (2.5)$$

The inequalities (2.4) and (2.5) imply

$$\begin{aligned} \frac{\varepsilon}{2} \|u'_\varepsilon(t+h) - u'_\varepsilon(t)\|^2 \leq \frac{1}{\varepsilon} \|u_\varepsilon(h) - u_\varepsilon(0)\|^2 + \frac{2}{\varepsilon} \left(\int_0^h \|Au_\varepsilon(s)\| ds \right)^2 + \frac{2}{\varepsilon} \left(\int_0^h \|f(s)\| ds \right)^2 + \\ + \int_0^t \|u'_\varepsilon(s+h) - u'_\varepsilon(s)\|^2 ds + \frac{1}{2} (A(u_\varepsilon(t+h) - u_\varepsilon(t)), u_\varepsilon(t+h) - u_\varepsilon(t)) + \\ + \int_0^t (f(s+h) - f(s), u'_\varepsilon(s+h) - u'_\varepsilon(s)) ds. \end{aligned} \quad (2.6)$$

We divide by h^2 and let $h \rightarrow 0$. In view of

$$\lim_{h \rightarrow 0} \frac{1}{h^2} \left(\int_0^h \|Au_\varepsilon(s)\| ds \right)^2 = \|Au_\varepsilon(0)\|^2$$

and

$$\lim_{h \rightarrow 0} \frac{1}{h^2} \left(\int_0^h \|f(s)\| ds \right)^2 = \|f(0)\|^2,$$

we see that (2.1) holds, so $u''_\varepsilon \in L^\infty(0, T; H)$.

3. THE CONSTRUCTION OF THE ASYMPTOTIC EXPANSION OF u_ϵ

Let us give now some formal definitions. Consider the equations

$$L_0 u = f_0, \tag{P_0}$$

$$L_0 u_\epsilon + \epsilon L_1 u_\epsilon = f_0 + \epsilon f_1, \tag{P_\epsilon}$$

where L_0, L_1 are two given operators, f_0, f_1 are known functions, $\epsilon > 0$ is a small parameter and $u(x), u_\epsilon(x), x \in D$ are the solutions of these problems, satisfying some initial or boundary conditions.

DEFINITION 3.1. We say that (P_ϵ) is regularly perturbed in the domain D if

$$\sup_{x \in D} \|u_\epsilon(x) - u_0(x)\| \rightarrow 0, \text{ as } \epsilon \rightarrow 0. \tag{3.1}$$

Otherwise, (P_ϵ) is called singularly perturbed problem.

Remark 3.1. We see that, in the case of the regularly perturbed problem, the solutions u and u_ϵ of $(P_0), (P_\epsilon)$ are close to each other and in the case of singular perturbations, u will not be in the proximity of (u_ϵ) , at least on a subset D_0 of D . This subset D_0 is called a boundary layer.

Let D_1 be a subset of D and $U_\epsilon(x)$ a function which is defined on D_1 . In particular, D_1 may coincide with D .

DEFINITION 3.2. The function $U_\epsilon(x)$ is said to be an asymptotic approximation for the solution u_ϵ of (P_ϵ) , with respect to the parameter ϵ , in the subdomain D_1 , if

$$\sup_{x \in D_1} \|u_\epsilon(x) - U_\epsilon(x)\| \rightarrow 0, \text{ as } \epsilon \rightarrow 0. \tag{3.2}$$

Furthermore, if

$$\sup_{x \in D_1} \|u_\epsilon(x) - U_\epsilon(x)\| = O(\epsilon^k), \tag{3.2}'$$

then we say that $U_\epsilon(x)$ is an asymptotic approximation for $u_\epsilon(x)$ in D_1 , to within accuracy of the order ϵ^k .

The symbol $\alpha(\epsilon) = O(\epsilon^k)$ means that, there exists $c, \epsilon_0 > 0$, such that,

$$\|\alpha(\epsilon)\| \leq c\epsilon^k, \text{ for all } 0 \leq \epsilon \leq \epsilon_0.$$

In this section, we shall present a method for the construction of an asymptotic approximation for the solution u_ϵ of (A_ϵ) , in entire interval $[0, T]$. Namely, we shall consider the series

$$\sum_{k=0}^{\infty} \epsilon^k u_k(t), \tag{3.3}$$

where $u_k \in L^2(0, T; H)$. The n -th partial sum

$$U_{n\epsilon}(t) = \sum_{k=0}^n \epsilon^k u_k(t), \tag{3.4}$$

will give us the asymptotic approximation of u_ϵ with an accuracy of order $\epsilon^{n+1/2}$, as we shall see in section 4.

DEFINITION 3.3. The series (3.3) which satisfies the above condition is called the asymptotic expansion (asymptotic series) for u_ϵ in the domain D , as $\epsilon \rightarrow 0$.

When we refer to the asymptotic series, we often mean an algorithm which allows us to find the terms of the asymptotic series, until a certain order n . We say that the asymptotic series is constructed when such an algorithm is presented.

Now, we come back to our problems (A_0) and (A_ϵ) . It is natural to ask whether $u(t)$ is the asymptotic approximation for $u_\epsilon(t)$ as $\epsilon \rightarrow 0$. Observe that $u(t)$ will not in general satisfy the boundary condition $u(T) = b$, therefore the problem (A_ϵ) is singularly perturbed, i.e. at least in some vicinity of the final point $t = T$, the function $u(t)$ will not be close to $u_\epsilon(t)$. This vicinity $(\delta, T]$ is said to be the boundary layer. But does $u(t)$ approximate $u_\epsilon(t)$ outside this vicinity? We wish to construct an asymptotic approximation for $u_\epsilon(t)$, which is valid in the whole interval $[0, T]$ and we also want to find its order of accuracy.

To this end, we are looking for an asymptotic expansion for the solution u_ϵ of (A_ϵ) , in a form that is quite typical for singularly perturbed problems:

$$u_\epsilon(t) = \bar{u}_\epsilon(t) + \Pi(\tau), \quad \text{with } \tau = (T - t)/\epsilon, \tag{3.5}$$

where $\bar{u}_\epsilon(t)$ is the regular power series and $\Pi(\tau)$ is a boundary layer series in the neighbourhood $(\delta, T]$ of $t = T$, i.e.

$$\bar{u}_\epsilon(t) = u_0(t) + \epsilon u_1(t) + \dots + \epsilon^k u_k(t) + \dots, \tag{3.6}$$

$$\Pi(\tau) = \Pi_0(\tau) + \epsilon \Pi_1(\tau) + \dots + \epsilon^k \Pi_k(\tau) + \dots \tag{3.7}$$

So, the asymptotic expansion has the form

$$u_\epsilon(t) = \sum_{k=0}^{\infty} \epsilon^k (u_k(t) + \Pi_k(\tau)). \tag{3.8}$$

Here, $\Pi_k(\tau)$ should be boundary functions, which means that they approach zero as $\tau \rightarrow \infty$. Hence, we assume a priori that

$$\Pi_k(\infty) = 0, \quad n = 0, 1, 2, \dots \tag{3.9}$$

The functions $\Pi_k(\tau)$ are important only in a neighbourhood of $t = T$ (i.e. $\tau = 0$), namely in the boundary layer $(\delta, T]$.

Using (3.5) in the differential equation from (A_ϵ) , it follows via the linearity of A

$$\epsilon \bar{u}_\epsilon''(t) + \frac{1}{\epsilon^2} \Pi''(\tau) - \bar{u}_\epsilon'(t) + \frac{1}{\epsilon} \Pi'(\tau) = A \bar{u}_\epsilon(t) + A \Pi(\tau) + f(t), \quad \text{a.e. on } [0, T]. \tag{3.10}$$

Substituting the series (3.6) and (3.7) into (3.10) and equating the coefficients of like powers of ϵ , separately for the terms depending on t and on τ , respectively, we infer

$$\begin{aligned} u_0'(t) + Au_0(t) &= -f(t) \quad \text{a.e. } t \in [0, T] \\ u_1'(t) + Au_1(t) &= u_0''(t) \quad \text{a.e. } t \in [0, T] \\ &\dots \dots \dots \\ u_k'(t) + Au_k(t) &= u_{k-1}''(t) \quad \text{a.e. } t \in [0, T] \end{aligned}$$

and

$$\begin{aligned} \Pi_0''(\tau) + \Pi_0'(\tau) &= 0 \quad \text{a.e. } \tau \in [0, \infty) \\ \Pi_1''(\tau) + \Pi_1'(\tau) &= A \Pi_0(\tau) \quad \text{a.e. } \tau \in [0, \infty) \\ &\dots \dots \dots \\ \Pi_k''(\tau) + \Pi_k'(\tau) &= A \Pi_{k-1}(\tau) \quad \text{a.e. } \tau \in [0, \infty). \end{aligned}$$

Now, let us substitute the expansion (3.8) into boundary conditions of (A_ϵ) and take into

account (3.9), we obtain

$$\begin{aligned} u_0(0) = a, \quad u_k(0) = 0, \quad n \geq 1, \\ u_0(T) + \Pi_0(0) = b, \quad u_k(T) + \Pi_k(0) = 0, \quad k \geq 1. \end{aligned}$$

If these equalities hold, then the series (3.8) satisfies the boundary condition $u_\varepsilon(T) = b$, hence we can say that the role of the series (3.7) is to satisfy, together with the series (3.6), the imposed boundary condition $u_\varepsilon(T) = b$. Thus, the Π -functions compensate for the discrepancy introduced by the regular part of the asymptotic expansion to the boundary condition $u_\varepsilon(T) = b$.

Unless (3.9), we need to require on the boundary functions the condition

$$\Pi_k \in H^2(0, \infty; H), \quad k \geq 0, \tag{3.11}$$

i.e. $\Pi_k, \Pi'_k, \Pi''_k \in L^2(0, \infty; H)$. It is known that, if $P \in H^2(0, \infty; H)$, then

$$\lim_{t \rightarrow \infty} \|P(t)\| = 0, \quad \lim_{t \rightarrow \infty} \|P'(t)\| = 0,$$

so (3.11) is stronger than (3.9).

Now, we can write the problems for the terms of the series (3.8). For the leading term $u_0(t)$ of the regular part of the asymptotic expansion, we obtain the boundary value problem

$$\begin{cases} u'_0(t) + Au_0(t) = -f(t), \quad \text{n.e. on } [0, T] \\ u_0(0) = a, \end{cases} \tag{3.12}$$

which evidently coincides with the unperturbed problem (A_0) .

For $u_k(t)$, $k \geq 1$, we deduce a problem of the same form with (3.12):

$$\begin{cases} u'_k(t) + Au_k(t) = u_{k-1}(t) \quad \text{n.e. on } [0, T] \\ u_k(0) = 0. \end{cases} \tag{3.13}$$

In the zeroth order approximation of the boundary layer series, we have

$$\begin{cases} \Pi''_0(\tau) + \Pi'_0(\tau) = 0 \quad \text{n.e. on } [0, \infty) \\ \Pi_0(0) = b - u_0(T), \quad \Pi_0 \in H^2(0, \infty; H) \end{cases} \tag{3.14}$$

and for the higher-order terms $\Pi_k(\tau)$, $k \geq 1$, of the series (3.7), we arrive at

$$\begin{cases} \Pi''_k(\tau) + \Pi'_k(\tau) = A\Pi_{k-1}(\tau) \quad \text{n.e. on } [0, \infty) \\ \Pi_k(0) = -u_k(T), \quad \Pi_k \in H^2(0, \infty; H). \end{cases} \tag{3.15}$$

Replace the hypothesis (I3) by the more restrictive condition

(I3)' for some fixed $f \in W^{n,2}(0, T; H)$ i.e. $f, f', \dots, f^{(n)} \in L^2(0, T; H)$.

$n \geq 1$.

(I4) $\alpha_k = -f^{(k-1)}(0) + Af^{(k-2)}(0) - \dots + (-1)^k A^{k-1} f(0) + (-1)^k A^k a \in D(A)$, $k = \overline{1, n}$.

Let us show that each of these problems has a unique solution (under the assumptions (I1), (I2), (I3)', (I4)).

First, (3.12) has a unique solution $u_0 \approx u \in W^{1,\infty}(0, T; H)$, as we explained in Introduction.

Remark 3.2. Let $w \in W^{1,2}(0, T; H)$ be the solution of

$$\begin{cases} w' + Aw = -f^{(n)}, \quad \text{n.e. } [0, T] \\ w(0) = \alpha_n. \end{cases}$$

Using the constants variation formula, we can easily see that $u^{(n)} = w$, so $u \in W^{n+1,2}(0, T; H)$.

Let us show now the uniqueness of the solution for (3.14) and (3.15).

PROPOSITION 3.1. *If (3.14) (or (3.15)) has solution $\Pi_k \in H^2(0, \infty; H)$, then it is unique.*

Proof. Let $\Pi_k, Q_k \in H^2(0, \infty; H)$ be two solutions of (3.14) (or (3.15)). Denoting by P the difference $P = P_k = \Pi_k - Q_k$, we observe that P verifies the problem

$$\begin{cases} P''(\tau) + P'(\tau) = 0, & \text{a.e. on } [0, \infty) \\ P(0) = 0, & P \in H^2(0, \infty; H). \end{cases} \quad (3.16)$$

Let τ be a given number in $[0, \infty)$. Multiplying (3.16) by P and integrating over $[\tau, \infty)$, we have

$$\int_{\tau}^{\infty} [(P'', P) + (P', P)] ds = 0$$

or, equivalently

$$(P', P)|_{\tau}^{\infty} + \frac{1}{2} \|P\|^2|_{\tau}^{\infty} = \int_{\tau}^{\infty} \|P'\|^2 ds. \quad (3.17)$$

Using the fact that $(P', P)(\infty) = 0$ and $\|P(\infty)\| = 0$, it follows that

$$(P'(\tau), P(\tau)) = -\frac{1}{2} \|P(\tau)\|^2 - \int_{\tau}^{\infty} \|P'\|^2 ds \leq 0, \quad \tau \geq 0 \quad (3.18)$$

so, the function $\tau \rightarrow (1/2)\|P(\tau)\|^2$ is nonincreasing on $[0, \infty)$. This implies that $P(\tau) = 0$, for every $\tau \in [0, \infty)$, hence we have the uniqueness.

Now, we want to solve the problem (3.14). To do this, we shall prove that there is $x \in H$ such that $\Pi_0(\tau) = \phi(\tau)x$ verifies (3.14), where $\phi : [0, \infty) \rightarrow \mathbb{R}$ is the solution of the boundary value problem

$$\begin{cases} \phi''(\tau) + \phi'(\tau) = 0, & \tau \in [0, \infty) \\ \phi(0) = 1, & \phi(\infty) = 0. \end{cases} \quad (3.19)$$

The solution of (3.19) is $\phi(\tau) = e^{-\tau}$, so $\Pi_0(\tau) = e^{-\tau}x$. From $\Pi_0(0) = b - u_0(T)$, we find

$$\Pi_0(\tau) = (b - u_0(T))e^{-\tau}, \quad \tau \in [0, \infty). \quad (3.20)$$

Remark that $\|\Pi_0(\tau)\| \leq ce^{-\tau}$, $(\forall)\tau \geq 0$, where $c = \|b - u_0(T)\| \geq 0$. The condition $\Pi_0 \in H^2(0, \infty; H)$ is also satisfied, because

$$\|\Pi_0\|_{L^2(0, \infty; H)}^2 = \|\Pi_0'\|_{L^2(0, \infty; H)}^2 = \|\Pi_0''\|_{L^2(0, \infty; H)}^2 = \frac{1}{2} \|b - u_0(T)\|^2 < \infty.$$

Therefore, we can state

PROPOSITION 3.2. *The unique solution Π_0 of (3.14) is given by (3.20).*

Remark 3.3. Similarly, we can find the unique solution of (3.15) for $k = 1$, namely,

$$\Pi_1(\tau) = -e^{-\tau}u_1(T) - \tau e^{-\tau}A(b - u_0(T)). \quad (3.21)$$

For an arbitrary $k \in \mathbb{N}$, we cannot find the exact solution of (3.15), but we can apply, for example, Theorem 1.1 to deduce that (3.15) has a solution $\Pi_k \in H^2(0, \infty; H)$. Indeed, we put $\tilde{r}(\tau) = e^{\tau}$, via (1.6), $B = 0$ (so $B_{\lambda} = 0$), α as in (1.9), $p = r = 1$ and $g_k = A\Pi_{k-1}$. In view of the mathematical induction method, if $\Pi_{k-1} \in L^2_{\tilde{r}}(0, \infty; H)$, it follows that $g_k \in L^2_{\tilde{r}}(0, \infty; H)$ and, according to Theorem 1.1, we find $\Pi_k \in H^2_{\tilde{r}}(0, \infty; H) \subseteq H^2(0, \infty; H)$.

So, we deduced that the problems (3.12) - (3.15) have unique solutions $u_k \in W^{2,2}(0, T; H)$, $\Pi_k \in H^2(0, \infty; H)$, $k = \overline{0, n}$. Thus, the algorithm described allows one to define the terms

of the series (3.6) and (3.7) up to the order n , inclusive (under the hypotheses (I1), (I2), (I3)').

Remark 3.4. If $f \in W^{1,2}(0, T; H)$ (i.e. we have (I3), not (I3)'), we can find only the zeroth order terms of the series, u_0 and Π_0 .

4. THE JUSTIFICATION OF THE ASYMPTOTIC EXPANSION

Let us denote by $S_{n\varepsilon}$ the n -th partial sum of the series (3.8), and by r_ε the remainder term, i.e.

$$S_{n\varepsilon}(t) = \sum_{k=0}^n \varepsilon^k [u_k(t) + \Pi_k(\tau)], \tag{4.1}$$

$$r_\varepsilon(t) = u_\varepsilon(t) - S_{n\varepsilon}(t). \tag{4.2}$$

We desire to estimate $\sup_{0 \leq t \leq T} \|r_\varepsilon(t)\|$. The remainder term verifies the boundary value problem (which has a similar form to (A $_\varepsilon$)):

$$\begin{cases} \varepsilon r_\varepsilon'' - r_\varepsilon' = Ar_\varepsilon + \varepsilon^n (A\Pi_n(\tau) - \varepsilon u_n'') \text{ a.e. on } [0, T] \\ r_\varepsilon(0) = r_\varepsilon(T) = 0. \end{cases} \tag{4.3}$$

Remark 4.1. First, observe that, for a given $n \in \mathbb{N}$, $u_n'' \in L^2(0, T; H)$, via Remark 3.2. Then, the solution Π_n of (3.15) or (3.14) is in $L^2(0, \infty; H)$. Moreover, since $\tau = (T - t)/\varepsilon$, then

$$\int_0^T \|A\Pi_n(\tau)\|^2 dt = \varepsilon^2 \int_0^\infty \|A\Pi_n(\tau)\|^2 d\tau < \infty.$$

Then, there exists $M_n > 0$ such that

$$\int_0^T \|A\Pi_n\left(\frac{T-t}{\varepsilon}\right) - \varepsilon u_n''(t)\|^2 dt \leq \varepsilon^2 M_n^2, \tag{4.4}$$

i.e. $A\Pi_n(\tau) - \varepsilon u_n''(t) \in L^2(0, T; H)$. Therefore, the problem (4.3) has a unique solution $r_\varepsilon \in W^{2,2}(0, T; H)$.

Now we can state the main result of this section, which will be the answer of the questions from Introduction.

THEOREM 4.1. *Under the hypotheses (I1), (I2), (I3)', (I4), for sufficiently small ε , the series (3.8) is the asymptotic expansion (series) of the solution u_ε of (A $_\varepsilon$) in the interval $[0, T]$, as $\varepsilon \rightarrow 0$, i.e. the following estimate holds*

$$\sup_{0 \leq t \leq T} \|u_\varepsilon(t) - S_{n\varepsilon}(t)\| = O(\varepsilon^{n+1/2}). \tag{4.5}$$

Proof. Without loss of generality, suppose that $0 = A0$.

Multiplying the equation of (4.3) by r_ε and integrating from 0 to T (for a fixed $n \in \mathbb{N}$), we deduce, via (4.4)

$$\int_0^T (Ar_\varepsilon, r_\varepsilon) dt + \frac{1}{2} \int_0^T (\|r_\varepsilon\|')^2 dt \leq \varepsilon \int_0^T (r_\varepsilon'', r_\varepsilon) dt + \varepsilon^{n+1} \cdot M_n \left(\int_0^T \|r_\varepsilon\|^2 dt \right)^{1/2}$$

Denote by $|\cdot|$ the norm in $L^2(0, T; H)$. Since A is positive definite, i.e. (1.1), then, the

above inequality and the boundary conditions $r_\varepsilon(0) = r_\varepsilon(T) = 0$ imply

$$\omega |r_\varepsilon| \leq \varepsilon |r_\varepsilon''| + \varepsilon^{n+1} \cdot M_n. \quad (4.6)$$

Next, one multiplies the equation in (4.3) by Ar_ε and one integrates over $[0, T]$. But A is symmetric, so, integrating by parts, we have

$$\int_0^T (Ar_\varepsilon, r_\varepsilon') dt = 0 \quad (4.7)$$

and

$$\int_0^T (Ar_\varepsilon, r_\varepsilon'') dt \leq -\omega \int_0^T \|r_\varepsilon'\|^2 dt \leq 0. \quad (4.8)$$

Then,

$$\int_0^T \|Ar_\varepsilon\|^2 dt = \varepsilon \int_0^T (r_\varepsilon'', Ar_\varepsilon) dt - \varepsilon^n \int_0^T (A\Pi_n(\tau) - \varepsilon u_n''(t), Ar_\varepsilon(t)) dt.$$

hence (4.8) and (1.1) give us

$$|Ar_\varepsilon|^2 + \omega \varepsilon |r_\varepsilon'|^2 \leq \varepsilon^{n+1} M_n |Ar_\varepsilon|. \quad (4.9)$$

From this inequality, we infer

$$|Ar_\varepsilon| \leq \varepsilon^{n+1} M_n \quad (4.10)$$

and then, coming back to (4.9), one arrives at

$$|r_\varepsilon'|^2 \leq \varepsilon^{2n+1} M_n^2 / \omega. \quad (4.11)$$

Let us multiply the equation from (4.3) by r_ε'' and let us integrate over $[0, T]$, to obtain, with the aid of (4.8) and (4.4),

$$\varepsilon |r_\varepsilon''| \leq |r_\varepsilon'| + \varepsilon^{n+1} M_n.$$

According to (4.11), this leads to

$$\varepsilon^2 |r_\varepsilon''|^2 < 2M_n^2 \varepsilon^{2n+1} (1/\omega + \varepsilon), \quad (4.12)$$

therefore (4.6) implies

$$\omega |r_\varepsilon| \leq K_1 \varepsilon^{(2n+1)/2} + \varepsilon^{n+1} M_n, \quad (4.13)$$

where K_1 is a positive constant, independent of ε .

But, $\|r_\varepsilon(t)\|^2 = 2 \int_0^t (r_\varepsilon(s), r_\varepsilon'(s)) ds \leq 2 |r_\varepsilon| |r_\varepsilon'|$. Then, in view of (4.11) and (4.13), it follows that, there is a positive K_2 , independent of ε , such that

$$\|r_\varepsilon(t)\|^2 \leq K_2 \varepsilon^{2n+1}, \quad t \in [0, T], \quad (4.14)$$

which leads to (4.5), as claimed.

Remark 4.2. Observe that $|Ar_\varepsilon| = O(\varepsilon^{n+1})$, $|r_\varepsilon'| = O(\varepsilon^{n+1/2})$, $|r_\varepsilon| = O(\varepsilon^{n+1/2})$.

Remark 4.3. For $n = 0$, we have $S_{0\varepsilon}(t) = u(t) + \Pi_0(\tau)$, hence, taking into account (3.20), we get

$$\|u_\varepsilon(t) - u(t)\| \leq \|u_\varepsilon(t) - S_{0\varepsilon}(t)\| + \|\Pi_0(\tau)\| \leq K_3 \sqrt{\varepsilon} + \|b - u(T)\| e^{(t-T)/\varepsilon} \quad (4.15)$$

with $K_3 > 0$ independent of ε . It is clear that, for every $t \in [0, T]$,

$$u_\varepsilon(t) \rightarrow u(t), \quad \text{as } \varepsilon \rightarrow 0, \quad (4.16)$$

but

$$\sup_{t \in [0, T]} \|u_\varepsilon(t) - u(t)\| \leq K_3\sqrt{\varepsilon} + \|b - u(T)\| \rightarrow \|b - u(T)\|, \text{ as } \varepsilon \rightarrow 0. \quad (4.17)$$

Evidently, if $b = u(T)$, this supremum has the limit zero and it is of order $\sqrt{\varepsilon}$; therefore, the problem (A_ε) is regularly perturbed in $[0, T]$. If $b \neq u(T)$, then u_ε approaches u just outside of a neighbourhood $(\delta, T]$ of T (for some δ , $0 < \delta < T$), because

$$\sup_{t \in [0, \delta]} \|u_\varepsilon(t) - u(t)\| \leq K_3\sqrt{\varepsilon} + \|b - u(T)\| \cdot e^{(\delta-T)/\varepsilon} \rightarrow 0, \text{ as } \varepsilon \rightarrow 0. \quad (4.18)$$

In this case, (A_ε) is singularly perturbed.

5. APPLICATION TO PARTIAL DIFFERENTIAL EQUATIONS

We are given two real Hilbert spaces V and H , such that $V \subseteq H$ and the inclusion mapping of V into H is continuous and densely defined. The space H is identified with its own dual H' and V' denotes the dual space of V . Then, the following inclusion relation holds:

$$V \subset H \subset V'.$$

It is known that, if A is a maximal monotone subset in $V \times V'$, then A_H defined by

$$\begin{aligned} D(A_H) &= \{v \in V; Av \in H\} \\ A_H v &= Av, \quad (\forall)v \in D(A_H), \end{aligned}$$

is a maximal monotone subset of $H \times H$.

We take $V = H_0^1(\Omega)$, $H = L^2(\Omega)$, $V' = H^{-1}(\Omega)$, where Ω is a bounded domain of \mathbb{R}^N with smooth boundary $\Gamma = \partial\Omega$. Let A be the linear operator given by

$$Au = -\Delta u = -\sum_{i=1}^N \frac{\partial^2 u}{\partial x_i^2}, \quad (5.1)$$

$$D(A) = H^2(\Omega) \cap H_0^1(\Omega). \quad (5.2)$$

Let a, b be two functions belonging to $D(A)$ and $f \in W^{n+1,2}(0, T; L^2(\Omega))$. Then the problems (A_0) and (A_ε) become

$$\begin{cases} \frac{\partial u}{\partial t} - \Delta_x u = -f, & \text{a.e. on } [0, T] \times \Omega \\ u = 0 & \text{a.e. on } [0, T] \times \Gamma \\ u(0, \cdot) = a(\cdot) \end{cases} \quad (5.3)$$

and

$$\begin{cases} \varepsilon \frac{\partial^2 u_\varepsilon}{\partial t^2} - \frac{\partial u_\varepsilon}{\partial t} = -\Delta_x u_\varepsilon + f, & \text{a.e. on } [0, T] \times \Omega \\ u_\varepsilon = 0 & \text{a.e. on } [0, T] \times \Gamma \\ u_\varepsilon(0, \cdot) = a(\cdot), \quad u_\varepsilon(T, \cdot) = b(\cdot), \end{cases} \quad (5.4)$$

respectively. For a fixed $n \in \mathbb{N}$, let $S_{n\varepsilon}(t, x)$ be the n -th partial sum of the series (3.8) and $r_\varepsilon = u_\varepsilon - S_{n\varepsilon}$. Then, for $a, b \in H_0^1(\Omega) \cap H^2(\Omega)$, $f \in W^{n+1,2}(0, T; L^2(\Omega))$ and sufficiently small ε ,

$$\|r_\varepsilon\|_{L^2([0, T] \times \Omega)} = O(\varepsilon^{n+1/2}) \quad (5.5)$$

and

$$\left\| \frac{\partial r_\epsilon}{\partial t} \right\|_{L^2([0,T] \times \Omega)} = O(\epsilon^n). \quad (5.6)$$

Here, r_ϵ verifies the linear problem

$$\begin{cases} \epsilon \frac{\partial^2 r_\epsilon}{\partial t^2} - \frac{\partial r_\epsilon}{\partial t} = -\Delta r_\epsilon - \epsilon^n \left(\Delta H_0 + \epsilon \frac{\partial^2 u}{\partial t^2} \right), & \text{a.e. on } [0, T] \times \Omega \\ r_\epsilon = 0 & \text{a.e. on } [0, T] \times \partial\Omega \\ r_\epsilon(0, x) = r_\epsilon(T, x) = 0 & \text{a.e. on } \Omega. \end{cases} \quad (5.7)$$

Since $\int_\Omega r_\epsilon^2 dx = 2 \int_0^t \int_\Omega r_\epsilon \frac{\partial r_\epsilon}{\partial t} dt dx$ a.e. $t \in [0, T]$, it follows that

$$\sup_{t \in [0, T]} \int_\Omega r_\epsilon^2(t, x) dx = O(\epsilon^{2n+1/2}).$$

In particular, for $n = 0$,

$$\sup_{t \in [0, T]} \left[\int_\Omega [u_\epsilon(t, x) - u(t, x)]^2 dx \right]^{1/2} \leq C \sqrt{\epsilon} + \int_\Omega [b(x) - u(T, x)]^2 dx.$$

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